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## Hydrodynamics of Manukau Harbour, New Zealand

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**Abstract** A calibrated two-dimensional finite-element model, which handles flooding/drying of intertidal areas, complemented by field data, has produced information on tidal propagation in Manukau Harbour, New Zealand with reasonably good accuracy. Excluding the upper-harbour inlets, c. 50% of both the amplification and phase lag in the dominant  $M_2$  tidal constituent occurs through the deep 10-km-long entrance channel to Puponga Point. The calculation of tidal residuals has isolated flood-directed and ebb-directed residual circulation respectively on either side of Puponga Point, which is a typical pattern for alternating flows around headlands. The model has established the dominance of tidal-driven over wind-driven circulation in most of the harbour except in the upper intertidal areas. Wind-driven circulation in

the inner Harbour is characterised by downwind flows over intertidal sand banks with pressure-driven return flows (vertically averaged) in the deeper main channels. Tidal dissipation rates are sufficiently high to inhibit the onset of any summer stratification.

**Keywords** tides; currents; numerical model; residual circulation; winds; estuary; Manukau Harbour

### INTRODUCTION

Manukau Harbour is the second largest coastal lagoon on the west coast of the North Island of New Zealand, adjacent to the city of Auckland (Fig. 1). It plays a major role in the coastal ecosystem because of its extensive intertidal sandbanks. It is extensively utilised as a fishery (pelagic and shellfish) and is the recipient of a major treated sewage discharge (mean discharge =  $3.9 \text{ m}^3 \text{ s}^{-1}$ ) from more than 600 000 Auckland residents and associated industries (Vant & Williams 1992).

Over recent years extensive studies of benthic community processes, phytoplankton dynamics, chemistry, and sediment transport within Manukau Harbour have been published (e.g., Pridmore et al. 1990; Vant 1991; Vant & Williams 1992; Thrush et al. 1994; Dolphin et al. 1995; Williamson et al. 1996; Thrush et al. 1997; Green et al. 1997). However, the most recent publication on the physical oceanography of Manukau Harbour, based on limited field measurements, was published 20 years ago (Heath et al. 1977). Increased knowledge of physical processes operating throughout the harbour is needed to elucidate their contributions to *in situ* biological, sedimentological, and chemical variability (both temporal and spatial). Of particular interest are circulation patterns, in both the channels and the extensive intertidal areas (40% of harbour), including the influence of wind. Two previous numerical modelling studies have been undertaken for Manukau Harbour, but with the prime objective to quantify the dispersion of effluent from potential



The surface area, which at mean high water spring tide (368 km<sup>2</sup>), is reduced by 60% at low spring tide, with the exposure of broad low-gradient tidal flats and banks covering 145 km<sup>2</sup>. Maximum depths across the entrance channel vary from 30 to 50 m in the deepest section off Paratutae Island. Inner-harbour channel depths vary from 30 m at Puponga Point to 5–10 m near the perimeter of the harbour. Minimum depths in the main shoaling sections are: Huia Banks in the entrance channel (5 m), Wairopa Channel (3 m), and Purakau Channel (2.5 m). The volume of the entrance channel at low water is sufficiently large to accommodate the entire tidal prism of the inner harbour (Heath et al. 1976; Vant & Williams 1992).

The catchment area (850 km<sup>2</sup>) is only about twice that of the harbour itself, so the mean-annual freshwater inflow (including sewage) of 26.3 m<sup>3</sup> s<sup>-1</sup> is relatively small (Vant & Williams 1992). This is reflected in the horizontal surface salinity distribution (measured in PSS), routinely monitored around high water, which shows a salinity **difference** between the entrance and the inland tidal inlets of only 4–5 over an average-rainfall year and up to 8 in a very wet season. The water in the harbour can be classified as well mixed, away from the influence of the sewage discharge from the Mangere Wastewater Treatment Plant, with an average salinity increase of only 0.2 down the water column (Vant & Williams 1992; W. N. Vant pers. comm.).

### Tides and winds

The tidal ranges in Manukau Harbour are among the highest in New Zealand (Hume et al. 1992). The mean spring (neap) tidal range is amplified inside the harbour, from 2.7 m (1.5 m) at the harbour entrance (Paratutae Island; Fig. 1) to 3.4 m (2.0 m) at Onehunga Wharf in Mangere Inlet (RNZN 1994). Measurements by Heath et al. (1977) showed that peak velocities of up to 1.8 m s<sup>-1</sup> can occur in the entrance channel. Other regions where high velocities have been measured are in the mouths of the inland tidal inlets: Mangere Inlet, 1.0 m s<sup>-1</sup> (spring) and 0.5 m s<sup>-1</sup> (neap) (Hume & Bell 1993; Green & Bell 1994); and Pahurehure Inlet, 0.9 m s<sup>-1</sup> (neap) (Hume 1979b). Mean tidal excursion distances for the whole harbour were succinctly summarised by harbour box-model segments (Vant & Williams 1992, fig. 1). These tidal excursions are confirmed by drogoue tracking data (Heath et al. 1977; Greig & Chiswell 1991). Estimates of the residence times for the upper reaches of the

Papakura Channel, Waiuku Channel, and the Wairopa/Purakau system, under mean annual hydrological conditions, are 12, 13, and 26 days respectively (Vant & Williams 1992). More recent work indicates that the residence time for Papakura Channel (i.e., entrance to Pahurehure Inlet), maybe similar to that of the northern channel system.

Prevailing surface-wind directions are from the south-west (average 26% of a year), north to north-east (24%), and west (10%), whereas it is calm for 13% of the time. The median wind speed is c. 6 m s<sup>-1</sup>. Sustained winds >15 m s<sup>-1</sup> are seldom experienced (0.2%), usually associated with cold, unstable south-west airstreams (New Zealand Meteorological Service 1982).

## METHODS

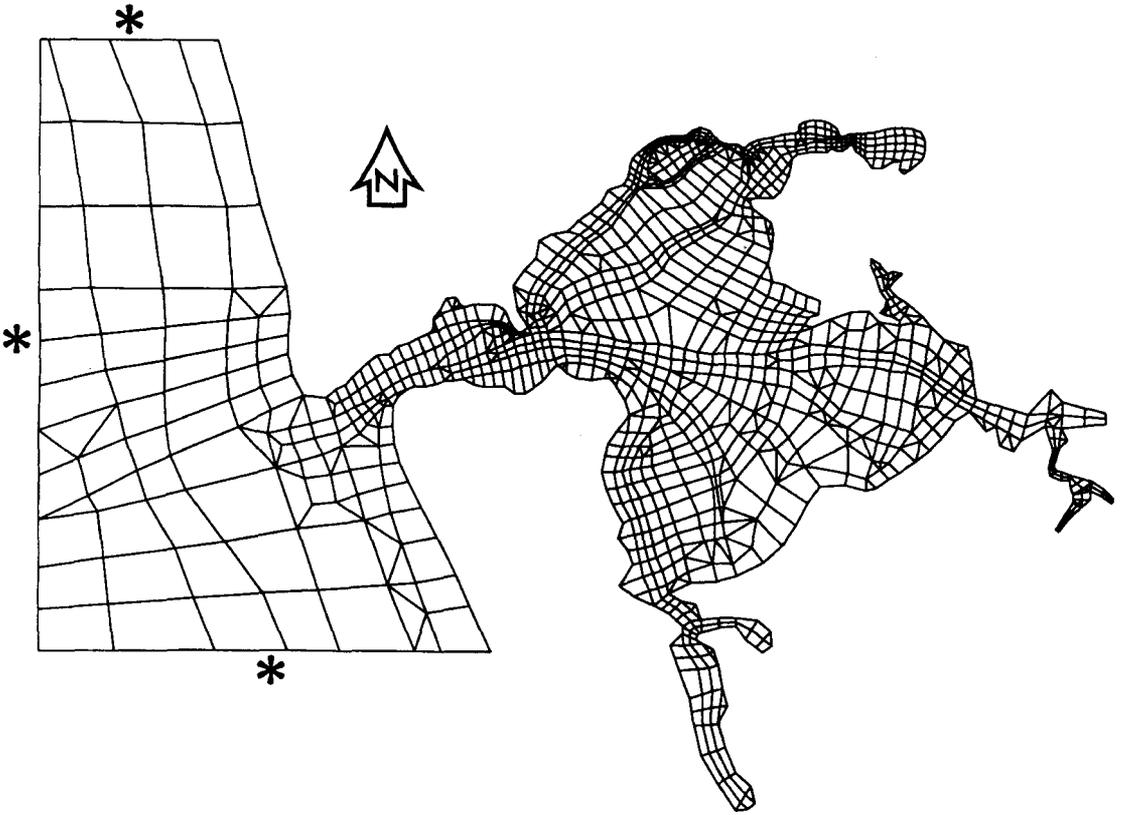
### Field studies

Field studies were concentrated during February–April 1994. Synoptic tidal heights were sampled over c. 1 month at several sites and current velocities at some of these (Fig. 1). Tide height and current measurements from other field studies were also utilised. A summary of instrumentation and deployment information for eight sites is given in Table 1.

Tidal height and current datasets were pre-processed by the removal of glitches and end sections effected by biofouling, then analysed for harmonic constituents using the algorithms of Foreman (1977, 1978). Simultaneous wind velocity and barometric pressure at hourly intervals were obtained from the climate station at Auckland Airport (Fig. 1) to assist with interpretation of water levels and velocities.

### Numerical tidal model

The finite-element model RMA-2 (King 1985, 1992; Pierson et al. 1995) was used to simulate the hydrodynamics of Manukau Harbour. This two-dimensional depth-averaged model solves the momentum and continuity equations, treating the two orthogonal velocity components ( $u$ ,  $v$ ) and water depth ( $h$ ) as the dependent variables (King 1992; Roig & King 1992). Boussinesq assumptions are incorporated as eddy viscosity terms to represent turbulent energy losses. Other terms built into the governing equations describe friction losses (approximated by a Chezy formulation), the Coriolis effect, and surface wind stresses. The



**Fig. 2** Finite-element mesh of Manukau Harbour, New Zealand. Asterisks denote where external sea level boundary conditions were applied.

The tide height throughout the model was initially set to a horizontal level at mean sea level (m.s.l.), which is 2.37 m above Chart Datum. A period of 2.5 tidal cycles was allowed for the model to equilibrate. A small model time step of 0.2 h (or

62 time-steps per  $M_2$  tide cycle) was used to ensure accurate system response to the input boundary conditions and minimise truncation errors. Wind effects were simulated by applying surface wind stresses for prevailing south-west and north-east winds, each for two speeds of 7.5 and 15  $\text{m s}^{-1}$ . The wind speed was increased gradually from zero to the required value over the equilibration period.

**Table 2** Sources of bathymetry used to generate the finite-element model mesh.

Harbour region	Bathymetry/shoreline source
Main harbour and coastal region	Chart NZ4314, RNZN (1984); Chart NZ4315, RNZN (1984)
Upper Mangere Inlet	Hume (1979a)
Pahurehure Inlet	NZMS 260, sheet R12; Hume (1979b)
Wiroa Island	Dolphin (1992); T. M. Hume (pers. comm.)
Waiuku River	T. M. Hume (pers. comm.)

Residual vectors of velocity, shear stress, and dissipation were integrated over a tidal cycle using the Romberg numerical integration technique. For surface-wind simulations, the vector **difference** in residual velocities (wind minus calm) was computed to isolate the component arising solely from the applied wind stress.  $M_2$  tidal amplitudes and phases were computed for all mesh nodes which remained submerged throughout the tidal cycle, using the method described by Foreman & Henry (1989).

## RESULTS AND DISCUSSION

### Water-surface elevations

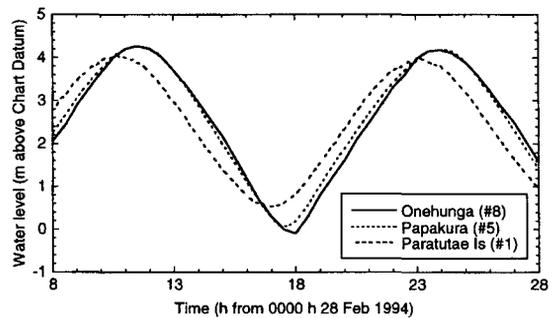
#### Tides

Harmonic analysis of the long-term tidal record from Onehunga Wharf (Site 8; Fig. 1) highlights the dominance of the semi-diurnal (Group 2) tidal constituents, mainly  $M_2$ ,  $S_2$ , and  $N_2$  in descending order (Table 3). These three semi-diurnal constituents account for 97.7% of the variance in the tidal signal compared to 0.2% for the main diurnal constituents ( $O_1$ ,  $K_1$ ). For Onehunga Wharf, the tidal form factor  $F$  is 0.045, which also classifies the characteristic tide as predominantly semi-diurnal (i.e.,  $F \ll 0.25$ ; Pugh 1987).

The mean perigean tide range  $2(M_2+N_2)$ , is slightly less at 3.24 m than the mean spring range,  $2(M_2+S_2)$ , of 3.42 m. This arises because the solar semi-diurnal tide ( $S_2$ ) is larger than the elliptical semi-diurnal tide ( $N_2$ ), in contrast to the east coast of New Zealand (Goring & Bell 1995; Bell & Goring 1996). Therefore the fortnightly spring/neap cycle (following the phases of the Moon) is more pronounced than the monthly perigean/apogean effect (synchronous with the Moon being in its perigee or apogee). Nevertheless with the  $N_2$  amplitude being 75% of the  $S_2$  amplitude, the latter effect is still important in Manukau Harbour, particularly when perigean tides coincide with

**Table 3** Main tidal constituents, by frequency groups, computed from harmonic analysis of the long-term tide gauge record for Onehunga Wharf (Site 8, Fig. 1) sourced from Ports of Auckland (unpubl. data), New Zealand. (Phases are relative to the tide-raising force at Greenwich for NZST zone.)

Constituent	Amplitude (cm)	Phase (°)
<b>Group 0 (long period)</b>		
$S_a$	9.2	26
<b>Group 1 (diurnal)</b>		
$O_1$	1.7	111
$K_1$	6.0	208
<b>Group 2 (semi-diurnal)</b>		
$N_2$	26.9	284
$M_2$	135.2	301
$S_2$	35.7	355
$K_2$	9.8	348
<b>Group 4 (fourth-diurnal)</b>		
$M_4$	5.5	43
<b>Group 6 (sixth-diurnal)</b>		
$M_6$	4.2	175



**Fig. 3** Comparison of measured spring tide height envelopes for Sites 1 (Paratutae Is.), 5 (Papakura Channel–Wiroa Is.), and 8 (Onehunga Wharf, New Zealand).

spring tides every 7.4 anomalistic months (203 days), causing higher than normal spring tides.

Because of the location of Onehunga Wharf in the upper reaches of the harbour, nearly 30 km from the entrance, the tidal record there also exhibits enhanced over-tides (Group 4 and above; Table 3). Over-tides arise in shallow waters as a result of bed-stress non-linearities. The main over-tides at Onehunga Wharf are  $M_4$  (6.2 h) and  $M_6$  (4.1 h) accounting for 0.3% variance. The remaining small variance in the water level record stems from other smaller amplitude Group 2 constituents (e.g.,  $K_2$ ) and low-frequency Group 0 components (Table 3) such as the non-tidal annual cycle ( $S_a$ ).

The  $M_2$ ,  $S_2$ , and  $N_2$  tidal amplitudes are amplified by 15–21% between the harbour entrance (#1) and the upper reaches of the main channels of the harbour. Phase (time) lags for these same constituents, relative to the harbour entrance, range from c. 40 min to nearly 1 h. In the northern sector, most of this retardation of the tide wave occurs seaward of Site 7 in Wairoa Channel, with only a further 4-min phase lag applying up to Onehunga Wharf (#8). Fig. 3 displays the difference in water levels throughout a spring tidal cycle between the entrance and two upper harbour sites (#5 and #8; Fig. 1). Because of the significant amplification in tidal amplitude between the entrance and Onehunga Wharf, the reduced level of LAT at Paratutae Island is c. 55 cm above LAT or Chart Datum at Onehunga Wharf (D. Hoskins, Ports of Auckland pers. comm.).

Although  $M_2$  dominates non-linear processes during tidal propagation through the harbour basin, its first harmonic (over-tide)  $M_4$ , can be used to

focus on tidal asymmetry in such estuarine systems and their morphological implications (Speer & Aubrey 1985). For the Manukau, the amplitude ratio ( $M_4/M_2$ ) increases about 4-fold from 0.013 at the entrance to about 0.05 at the upper end of Papakura and Wairoa Channels, with c. 60% of this increase occurring between the entrance and nearest inner-harbour Site 3 (Fig. 1). The relative phase ( $2M_2 - M_4$ ) lies in a narrow band, starting at  $170^\circ$  at the entrance (Paratutae Island), rising to a peak of  $200^\circ$  at Sites 2 and 3, then steadily decreases to  $180 - 190^\circ$  at the upper harbour Sites 5, 8, and 9. These results are discussed later in the context of current patterns.

For eight sites throughout the Harbour (Fig. 1), both observed and modelled  $M_2$  amplitudes ( $A_o, A_m$ ) and phases ( $g_o, g_m$ ) are compared in Table 4. The mean difference between modelled and observed amplitude was  $-1$  cm and for phase,  $2^\circ$ . An overall measure of the match between a modelled and observed harmonic constituent is given by the vectorial difference ( $\Delta D$ ) computed as distances in the complex plane (Foreman et al. 1995):

$$\Delta D = \sqrt{(A_o \cos g_o - A_m \cos g_m)^2 + (A_o \sin g_o - A_m \sin g_m)^2} \quad (1)$$

For all sites,  $\Delta D \leq 9.5$  cm, with the largest  $M_2$  differences in the upper reaches of Papakura Channel (Sites 4 and 5), whereas in the northern sector of the Manukau, the  $M_2$  tide was more closely matched by the model, with  $\Delta D < 4$  cm. The general spatial features of the dominant  $M_2$  tide are shown by the cotidal plots (Fig. 4). The contours have largely been interpolated from the harmonic analysis of model results, but incorporate relevant observed  $M_2$  values where consistent differences from the model occurred, such as the upper Papakura Channel (Table 4). (Note: contours are only approximate for the large intertidal banks.)

The  $M_2$  amplitude attenuates slightly through the harbour entrance (from 111 cm off the ebb-tide delta to 110 cm near Paratutae Island) as the tidal flow is constricted through the narrow entrance (Fig. 4A). Thereafter there is a rapid 10 cm amplification of the  $M_2$  constituent through the deep, narrow entrance channel until it bifurcates into the multiple inner-harbour channels at Puhanga Point. Throughout the main inner-harbour basin, another 10–12 cm amplification occurs. Within the southern tidal inlet (Waiuku River), where some observed tidal information is available, the  $M_2$  tide continues to amplify (and slow) moving upstream (Fig. 4) to reach an amplitude of c. 140 cm at the head of the

inlet (Tonkin & Taylor 1986). Cophase lines (Fig. 4B) show that about half the total phase lag of the  $M_2$  tidal wave within Manukau Harbour occurs through the confines of the entrance channel, followed by a more slowly varying retardation within the broad, shallower, inner Harbour. The flood-tide delta, Huia Banks (Fig. 1), appears to be the major control on the propagation of the  $M_2$  tidal wave within the entrance channel.

#### *Non-tidal effects*

The super-elevation of the mean tide level between the harbour entrance (Paratutae Island) and Onehunga Wharf observed in the  $M_2$  model run was 4.5 cm under calm conditions, which concurs with an estimate of 6 cm derived from survey measurements (Tonkin & Taylor 1986). Super-elevation is caused by the differential celerities of the tidal wave between low and high tides in shallow harbours, requiring a set-up of the half-tide level to redress the imbalance in flow capacity between the wave crest and trough (Ippen 1966; Druey et al. 1983). During persistent wind events, an additional water-surface gradient (set-up or set-down) is generated in Manukau Harbour. Predicted spatial changes in water level arising from either north-east or south-west winds were compared with calm conditions. The greatest differences (set-up or set-down) under windy conditions arise in the shallower inner harbour, rather than in the deeper entrance channel where tidal current velocities are considerably higher. At Onehunga Wharf, during  $15\text{-m s}^{-1}$  south-west winds, additional set-ups of 3–4 and 6 cm are predicted by the model for High Water (HW) and Low Water (LW) respectively, whereas for north-east winds at the same speed, set-downs of  $-3$  and  $-6$  cm respectively. The reason for the higher differences at mean LW is that wind-driven set-up varies at any location according to the water depth, with the wind-stress term in the momentum equation depending inversely on the depth of water. During the main field programme the maximum total water-level set-up occurred when a south-west wind peaked at a mean speed of  $20\text{ m s}^{-1}$  (21 April 1994). The maximum water level set-up was 45 cm above tidal predictions, comprising wind-stress set-up within the harbour (estimated to be 10+ cm) together with the inverted barometer effect (low atmospheric pressure) and externally-forced coastal set-up.

The annual cycle constituent ( $S_a$ ) is essentially non-tidal, with the main contributor being the thermo-steric effect on sea level from the seasonal

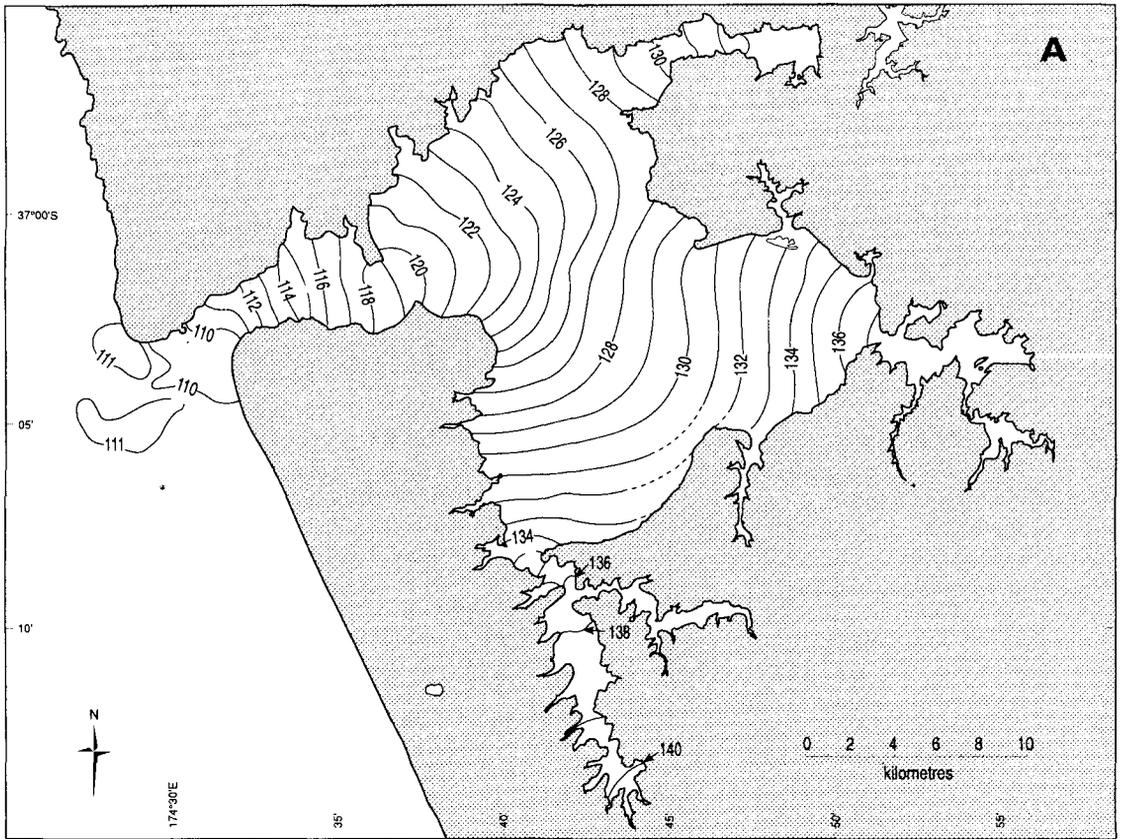


Fig. 4  $M_2$  cotidal diagrams: A, coamplitude (cm); and B, cophase (degrees).

heating/cooling cycle (Bell & Goring 1996). This is confirmed by the  $S_a$  phase of  $26^\circ$  (Table 3), with the peak height occurring at the end of January (mid summer). The resulting mean  $S_a$  amplitude of 9 cm

**Table 4** Comparison of observed and modelled  $M_2$  tide-elevation amplitudes ( $A_o$ ,  $A_m$ ) and phases ( $g_o$ ,  $g_m$ ) for numbered sites shown in Fig. 1 and the vectorial difference  $D$  (Equation 1).

Site	Observed		Model		$\Delta D$ (cm)
	$A_o$ (cm)	$g_o$ ( $^\circ$ )	$A_m$ (cm)	$g_m$ ( $^\circ$ )	
1	109.5	280.8	109.8	281.5	1.4
2	131.3	299.5	127.9	302.0	6.6
3	122.1	293.5	123.4	295.8	5.1
4	125.4	293.8	125.5	297.5	8.1
5	132.5	298.8	131.2	302.9	9.5
6	126.0	295.7	126.2	297.5	4.0
7	129.9	299.8	128.0	301.3	3.9
8	132.4	302.4	129.0	302.7	3.5

at Onehunga Wharf is somewhat higher than the amplitude of around 5–6 cm obtained from Geosat altimetry for the ocean off the north-west coast of New Zealand (Jacobs et al. 1992). The larger amplitude can be attributed to the greater range of seasonal density variability (water temperature and salinity) and summer evaporation losses within the shallower waters of Manukau Harbour (Vant & Williams 1992).

## Currents

### Tidal-current characteristics

From the harbour entrance up to the mid-reaches of the main inner-harbour channels (e.g., Sites 2, 5, 7; Fig. 1), the peak current velocities coincide with mid-tide on both flood and ebb. Slack tides throughout the harbour occur within 0–15 min after the respective local LW or HW, except the near-shore entrance region off South Head (Fig. 1), where slack tide lags local LW by c. 30 min. The overall

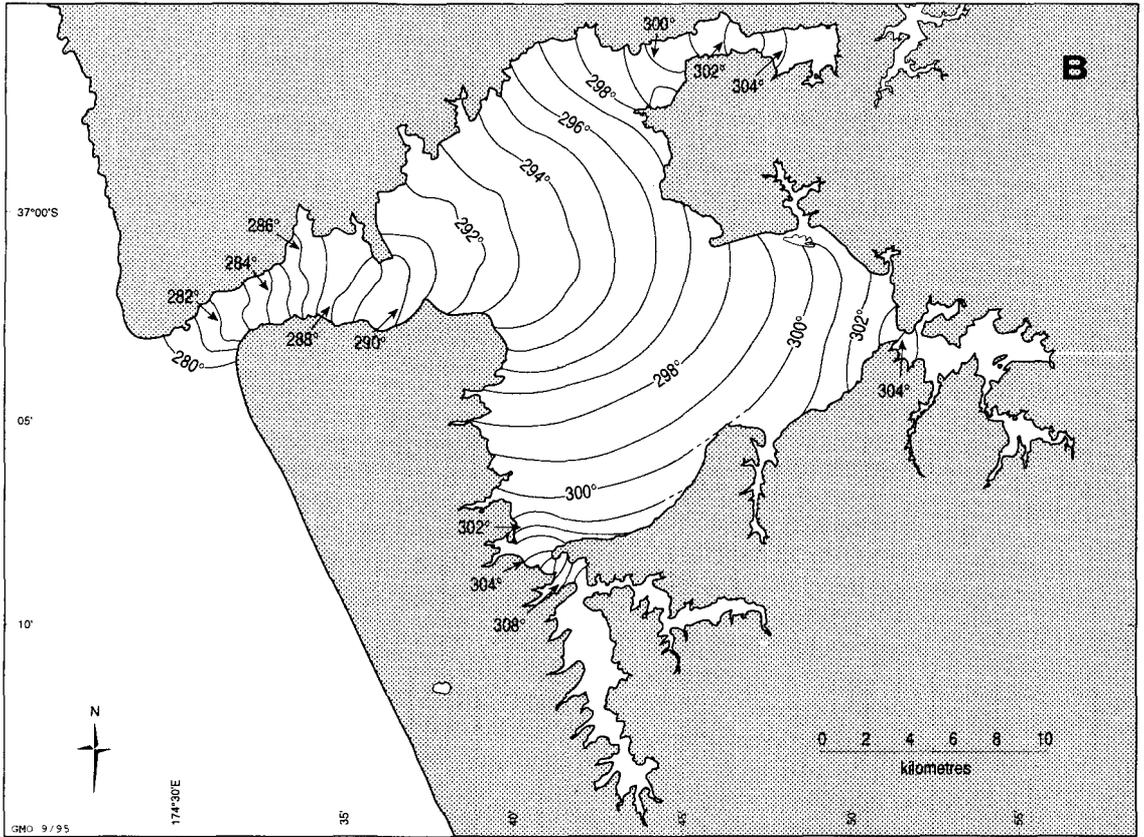


Fig. 4 (continued)

phase behaviour characterises the shallow-water tidal wave as a standing wave, as distinct from a progressive wave, and typically occurs in semi-enclosed basins where the tide wave is reflected back on itself (Ippen 1966; Druery et al. 1983). However towards the upper reaches of the harbour, there is a gradual distortion in the velocity phase, with peak currents occurring closer to HW. For example at Onehunga Wharf, the peak-flood velocity is only 1.75 h before HW whereas peak ebb occurs 2.5 h after HW (Hume & Bell 1993; Green & Bell 1994). The  $M_2$  tidal wave in the upper reaches of the Harbour is therefore not strictly a “standing” wave. However with the dominance of friction in such shallow-water systems, both the incident and reflected waves will rarely have similar magnitudes or properties, thus this type of wave classification is suspect (Jay 1991).

Tidal ellipse parameters for various non-overlapping current-meter records of at least 20

days for Papakura and Purakau Channels, are listed in Table 5. Because of the shortness of most records (<28 days), the  $M_2$  amplitudes obtained from harmonic analysis contain a contribution of c. 10% from the unresolved  $N_2$  constituent, which was established from the longer 30-day record at the lower meter in Papakura Channel (#4). Fitting a logarithmic velocity profile to the lower and upper current meter records at Site 4 to derive a depth-averaged  $M_2$  velocity at  $0.4 \times$  depth, and incorporating the 10% factor yields a depth-averaged amplitude (major semi-axis) of c.  $0.6 \text{ m s}^{-1}$  for Papakura Channel, whereas in the smaller Purakau Channel (#6) the depth-averaged  $M_2$  velocity was lower at  $0.48 \text{ m s}^{-1}$ . Adding  $S_2$  depth-averaged velocities of c.  $0.20$  and  $0.18 \text{ m s}^{-1}$  for respective channels yields depth-averaged mean spring-tide velocities of  $0.85 \text{ m s}^{-1}$  for Papakura Channel (#4) and  $0.7 \text{ m s}^{-1}$  for Purakau Channel (#6). As expected, the two dominant tidal current

constituents ( $M_2$  and  $S_2$ ) exhibited strongly rectilinear flow characteristics along the channel axis, shown by the near-zero ellipticity values (Table 5). Strong vertical shear is also evident in the velocity profile at the Papakura Channel site (#4) where the lower meter ( $0.14 \times$  depth) registered major-axis amplitudes ( $M_2$  and  $S_2$ ) which were 35% less than the upper meter ( $0.65 \times$  depth). The inclination of the major tidal axis was found to be nearly constant through the vertical for the dominant  $M_2$  and  $S_2$  components at Site 4, whereas the phases indicate that flow reversal near the bed occurs within a few minutes of that higher in the water column (Table 5).

#### Tidal circulation

The large tidal range generates strong  $M_2$  tidal currents, particularly in the entrance channel, as shown by the modelled velocity vectors gridded over the harbour (Fig. 5), for mid-flood and mid-ebb relative to the tide state off Puponga Point. Localised maximum ebb-tide velocities for a  $M_2$ -tide (spring-tide) reached  $1.9 \text{ m s}^{-1}$  ( $2.5 \text{ m s}^{-1}$ ) in the middle of the harbour entrance (Fig. 6B), opposite Ninepin Rock (see Fig. 1), where it begins to shoal towards the offshore delta and  $1.7 \text{ m s}^{-1}$  ( $2.2 \text{ m s}^{-1}$ ) in the southern entrance channel off Te Pirau Point (see Fig. 1), where the ebb-flow is locally constricted by the flood-tide delta (Huia Banks). These peak depth-averaged speeds are comparable to a previous peak spring-tide measurement of  $2.25 \text{ m s}^{-1}$  at the surface (Heath et

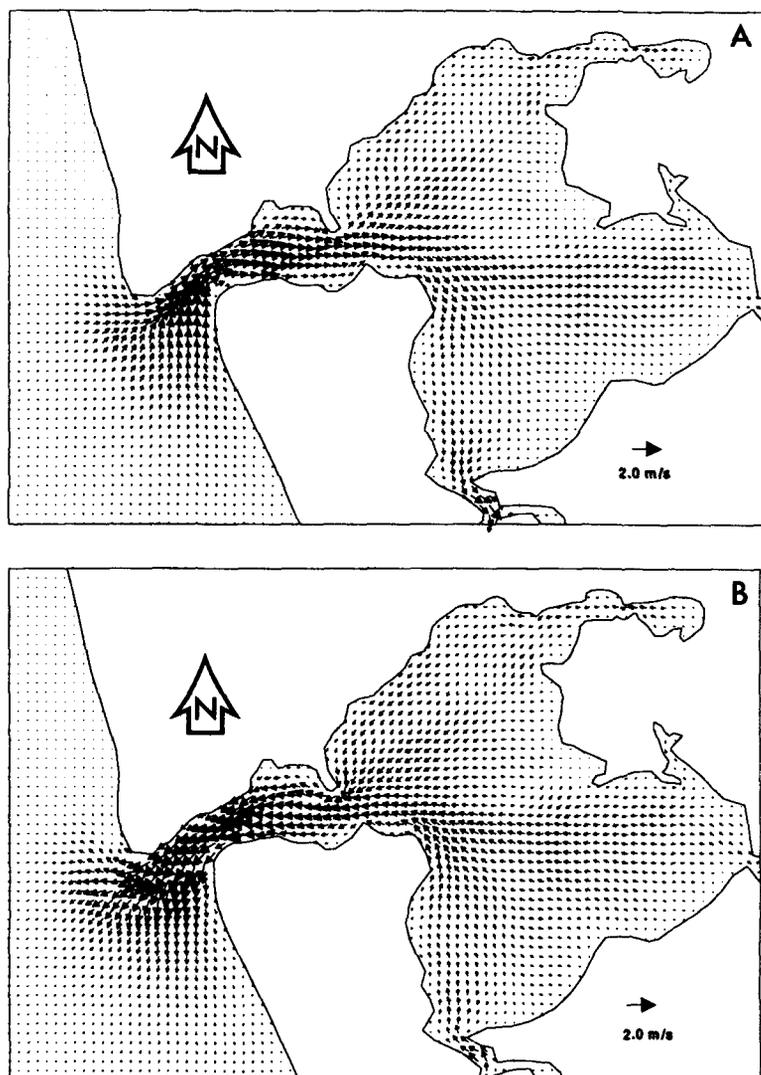
al. 1977). Maximum flood-tide velocities for  $M_2$ -tide (spring-tide) were somewhat lower, reaching  $1.5 \text{ m s}^{-1}$  ( $1.8 \text{ m s}^{-1}$ ) in both the nearshore flood-tide channel off South Head (Fig. 5A, 6D), and off Te Pirau Point (entrance channel). In a transect across the harbour entrance (Ninepin Rock–South Head), peak velocities during the ebb tide exceeded those for the flood tide, except for the southern 20% of the entrance width, where the nearshore flood-tide channel dominates. Further, at this southern end, the flood tide was 1 h longer than the ebb tide, otherwise the durations of both ebb and flood cycles were similar.

Although these results imply that for sediment transport the entrance is ebb-dominated with respect to peak velocities, in terms of the entire harbour, it is only partly consistent with a simplified sedimentary evolutionary sequence for enclosed basins discussed by Lincoln & Fitzgerald (1988) and Speer & Aubrey (1985). Basins such as the Manukau Harbour (a wide drowned river valley), which have been filled with low-lying intertidal banks, but incised with relatively deep channels ( $M_2$  amplitude/depth ratio  $\approx 0.1$ ), possess non-linear hydraulic filling characteristics (water surface area versus tidal elevation) which usually result in an ebb-dominated estuary. However the type of tidal distortion in estuaries (flood-dominant versus ebb-dominant) depends on the relative phasing of the overtide  $M_4$  to  $M_2$  (Speer & Aubrey 1985). Based on tidal elevations discussed earlier, Manukau Harbour exhibits relative phasing in the range

**Table 5** Tidal-current ellipse parameters for  $M_2$  and  $S_2$  from mid-channel current-meter deployments (20 days) in Papakura and Purakau Channels, New Zealand. (Ellipticity is the ratio of minor/major semi-axis, phase is relative to maximum ebb velocity, inclination is the orientation of the major axis relative to True North, and  $Z_0$  is the mean residual velocity.)

Site	Height/ water depth (m)	Constituent	Major semi- axis ( $\text{m s}^{-1}$ )	Minor semi- axis ( $\text{m s}^{-1}$ )	Ellipticity	Phase ( $^\circ$ )	Inclination ( $^\circ\text{T}$ )
<b>Papakura Channel</b>							
#4	4/28	$M_2$	0.53	0.01	0.02	37.8	279
		$S_2$	0.15	0.02	0.13	72.0	277
		$Z_0$	0.04				303 (ebb)
	18/28	$M_2$	0.82	0.01	0.01	36.0	278
		$S_2$	0.23	0.01	0.06	74.0	277
		$Z_0$	0.02				322 (ebb)
<b>Purakau Channel</b>							
#6	7/17	$M_2$	0.54	-0.01	-0.02	30.1	60
		$S_2$	0.18	-0.00	-0.01	73.5	63
		$Z_0$	0.05				69 (flood)

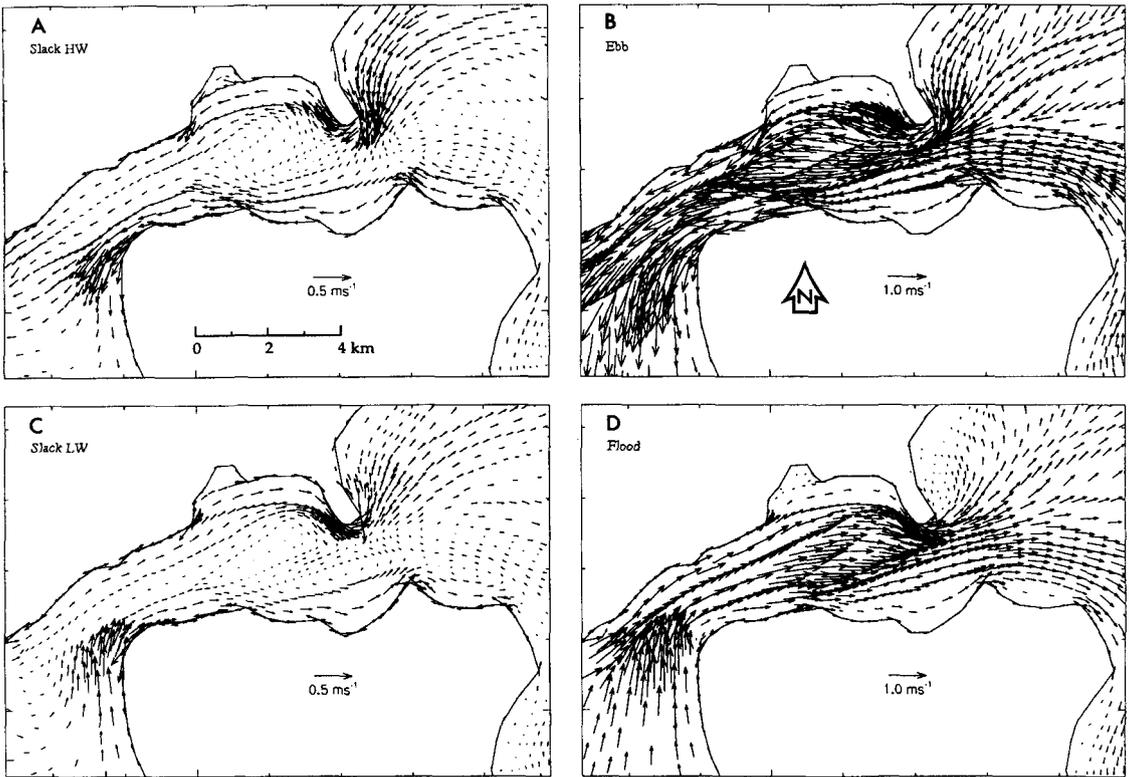
**Fig. 5** Modelled  $M_2$  current-vector fields for Manukau Harbour, New Zealand: **A**, peak-flood; and **B**, peak-ebb, relative to tide state in mid-channel off Puponga Point.



170–200°, which straddles the neutral value of 180°, thus not clearly favouring either an ebb-dominated or flood-dominated regime. This neutrality is probably because of the compensatory nature of residual flows in the multiple-branching channels within the Manukau basin, which forms a more complex circulation system than the single main channel assumed by Speer & Aubrey (1985).

The most complex velocity patterns in Manukau Harbour arise in the 9-km-long entrance channel and around Puponga Point headland (Fig. 6), which are strongly influenced by the bathymetry and complex shoreline geometry. Three significant

features moderate these complex current patterns: (1) the confluence of the major inner-harbour channels; (2) the constriction imposed by the Puponga Point headland; and (3) the flood-tide delta (Huia Banks) with a minimum water depth of 5 m, which forms a major obstacle to the flow in the middle of the deep entrance channel. A common feature at both HW and LW (Fig. 6A,C) is the early resumption of the next tide phase (flood and ebb respectively) along the nearshore areas on both sides of the entrance channel and the flood-tide channel off South Head in comparison with the mid-channel region and Huia Banks. At both HW and LW,



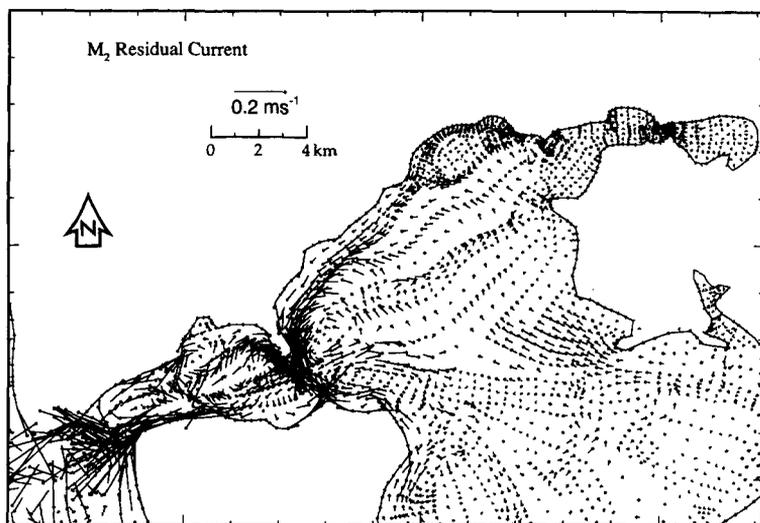
**Fig. 6** Modelled  $M_2$  current-vector fields for the Manukau Harbour entrance channel and Puoponga Point: **A**, slack high water (HW); **B**, peak-ebb; **C**, slack low water (LW); and **D**, peak-flood. (Note: change of velocity scale for slack-tide plots.)

velocities already have reached  $0.5 \text{ m s}^{-1}$  in the South Head channel. A small eddy is present between Mako and Kauri Points (Fig. 1), set up by the inertia of the decelerating flow in the large Papakura Channel and the nearshore tidal stream which changes direction earlier. For the peak-ebb tidal stream (Fig. 6B), the majority of streamlines across 75% of the entrance channel off Puoponga Point originate from the larger southern basin (Papakura/Waiuku Channels). The remaining streamlines (25% of flow width), which tightly converge around Puoponga Point on the northern side, link the northern basin channels to the entrance channel via the northern side of Huia Banks. This modelled behaviour is also confirmed by visual observations, which often show a sharp slick line relatively close to Puoponga Point, delineating waters of distinct water clarity originating from southern and northern sectors of the Harbour (W. N. Vant pers. comm.). At peak-flood (Fig. 6D), the flow

from the entrance channel north of Huia Banks is again tightly constrained around Puoponga Point, but flows mainly into the Purakau Channel. This circulation pattern leaves a weaker flood-tide flow in the northernmost Wairopa Channel, arising from a weak separation eddy in the lee of Puoponga Point, which enhances the ebb-dominance of the lower Wairopa Channel. The tightly constrained flow convergence around the Puoponga Point headland and, to a lesser extent, around Huia Banks, appears to curtail exchange of water and solutes between the northern and southern channels of the inner harbour as shown by dispersion modelling by Black et al. (1995).

For the inner harbour, peak  $M_2$ -tide (spring-tide) velocities  $> 0.6 \text{ m s}^{-1}$  ( $> 0.8 \text{ m s}^{-1}$ ) occur in the four main tidal channels (Fig. 5), particularly in the largest channel (Papakura) and at constricted entrances to the inland tidal inlets (Waiuku River, Mangere Inlet, and Pahurehure Inlet). The directions

**Fig. 7** Modelled  $M_2$  residual currents for Manukau Harbour, New Zealand.



of the tidal currents are generally tangential to the basin isobaths, but because bottom friction strongly affects current speeds, their magnitudes are proportional to mean water depth. Hence currents on the intertidal flats (Fig. 5) are markedly slower than adjacent channels accompanied by a greater tendency for velocities to be directed normal to the isobaths higher up on the flats as they drain or flood (Bell et al. 1997). No large-scale eddies in flow patterns are evident in the inner-harbour because of the relatively narrow, incised, channels. The distance out to the harbour entrance from the both Mangere Inlet (via Wairopa and Purakau Channels) and the outlet from Pahurehure Inlet (via the Papakura Channel), is equivalent to about three tidal excursions on a mean ebb tide (Vant & Williams 1992). Overall, the tidal circulation in Manukau Harbour is primarily governed by volume continuity, where the tidal currents diminish towards the harbour perimeter as both the tidal prism and depths decrease, except at localised inlet constrictions.

#### *Response to wind*

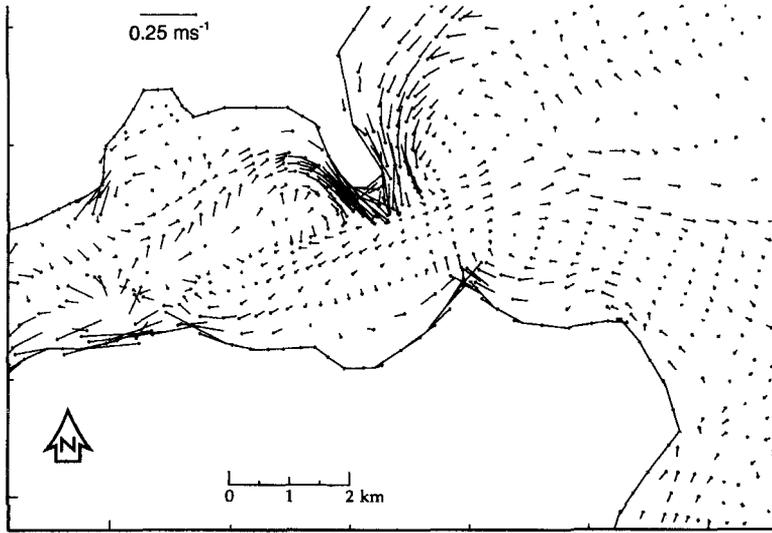
Tidal currents in the main channels are not appreciably modified because of either north-east or south-west winds of  $15 \text{ m s}^{-1}$ , with no discernible change in the deep entrance channel, where current velocities are substantial. The main effect of such wind events is to deviate the current velocities downwind along the shallow, exposed margins of the inner harbour, which is discussed further in the

next section. Tide and wind-generated currents in combination with locally generated wind waves, act together to suspend and advect fine sediments from the exposed intertidal banks, leading to enhanced turbidity levels which are a common feature of the water clarity in Manukau Harbour (Vant 1991; Dolphin et al. 1995; Green & Bell 1995; Bell et al. 1997; Green et al. 1997).

### **Residual currents**

#### *Tidal residuals*

Figure 7 shows the vector field of the Eulerian tidal residual currents over northern and central sectors of Manukau Harbour, and in Fig. 8, a closer view is presented for the region around Puponga Point. Tidal residual currents are generated by the interaction of tidal currents with bottom topography and harbour geometry and are manifest by ebb tides flowing for different periods than flood tides, differing ebb and flood current magnitudes, or by water masses following different paths. Effective residual mass or solute transport (which includes the covariation of tidal velocity and water depth during a tidal cycle) is enhanced in regions where the Eulerian residual-velocity field is spatially coherent over scales comparable to the tidal excursion. Coherent residual patterns on such scales within the inner-harbour only occur in the Wairopa Channel. Large-scale ebb-directed residual patterns occur in both the lower Wairopa Channel (with moderate residual velocities of up to  $0.1 \text{ m s}^{-1}$  towards Puponga Point) and along the southern



**Fig. 8** Modelled  $M_2$  residual current for the Manukau Harbour entrance channel and around Puponga Point headland, New Zealand.

shoreline of the entrance channel from Te Pirau Point through to the entrance (with strong residual velocities of up to  $0.35 \text{ m s}^{-1}$ ), where in both instances these areas are in the lee of the main flood tidal streams. Consistent flood-directed residuals occur along the northern flanks of Huia Banks, continuing on around the seaward side of Puponga Point headland (Fig. 8). The resultant pattern of Eulerian residual velocities generated either side of Puponga Point, and to a smaller extent, Mako Point (Fig. 8), is similar to the residual flow of counter-rotating phase eddies associated with strong oscillatory flows around headlands (Signell & Geyer 1990). However in this instance, development of the eddies is curtailed by the narrow channels and adjacent sandbanks around both sides of the headland compared to the open coast situation (Hume et al. 1997).

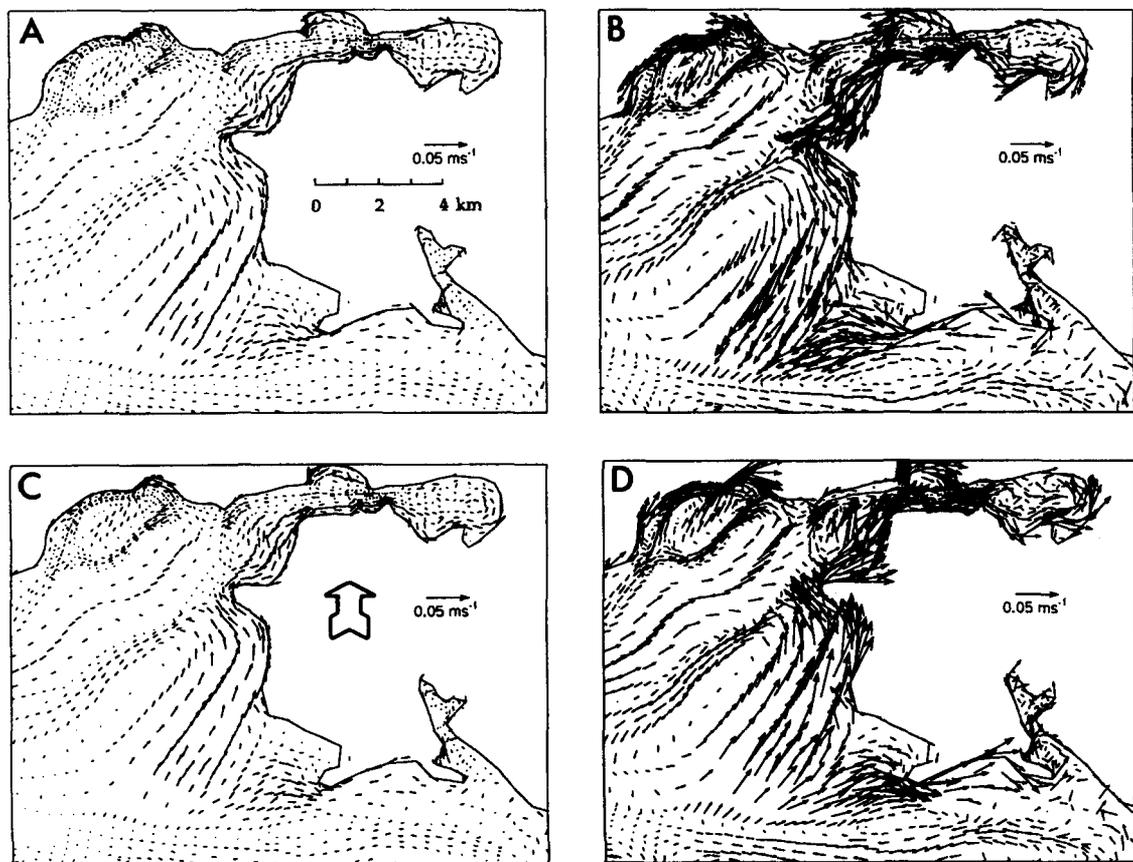
The residual flow pattern immediately seaward of Manukau Harbour entrance is typical of ebb-tide deltas (Oertel 1988), where sediments are expelled through an ebb-dominated flow through the central region of the delta, whereas a flood-dominated flow occurs along the nearshore zones of the coastal barriers. However as a consequence of the offset coastal barriers either side of the entrance, the flood tidal stream is strongest through the nearshore flood-channel off South Head (Fig. 6D).

#### *Steady-state wind response*

The response of residual currents in Manukau Harbour to steady-state winds was examined for

prevailing south-west and north-east winds of  $7.5$  and  $15 \text{ m s}^{-1}$  and compared with calm conditions. Vector differences in depth-averaged Eulerian residuals (wind minus calm) from the model results (Fig. 9), show the depth-averaged wind-induced circulation in the north-east sector of the Harbour. The wind-driven motion for the two north-east wind scenarios are shown in Fig. 9A,B and similarly for south-west winds in Fig. 9C,D. The first noticeable feature of the results is the non-linear response to increasing wind speed, where the wind-induced flows are markedly higher (more than  $2\times$ ) for  $15 \text{ m s}^{-1}$  winds, represented by the wind-stress term which is quadratic in wind speed. Second, the greatest wind response occurs on the shallow intertidal flats, rather than in the channels, arising from the inverse-depth relationship in the wind-stress term.

During **north-east** winds (Fig. 9A,B), consistent downwind flows occur along the entire north-eastern perimeter, commencing inside Mangere Inlet. A westerly-directed flow off Wiroa Island and past the Airport on the northern side of Papakura Channel merges with the residual flow along the north-eastern perimeter. Downwind wind-induced flows also occur along the entire length of Te Tau Bank and along the northern branch of Wairopa Channel, where it shoals. As discussed previously, the set-down in water elevations in the eastern sector of the harbour, arising from a north-east wind, creates water-surface gradients which drive residual vertically-integrated counter-flows upwind in the



**Fig. 9** Wind-induced residual circulation for the north-east sector of Manukau Harbour, New Zealand for—north-east winds: **A**,  $7.5 \text{ m s}^{-1}$ ; and **B**,  $15 \text{ m s}^{-1}$ ; and south-west winds: **C**,  $7.5 \text{ m s}^{-1}$ ; and **D**,  $15 \text{ m s}^{-1}$ .

main channels, as shown in Fig. 9B, although in the surface layer, residual flows may well be downwind. In particular, the total flood-directed residual flow in Purakau Channel (south-west of the Mangere WTP) is virtually doubled to  $0.05 \text{ m s}^{-1}$  with the addition of the wind-induced counter-flow from a  $15\text{-m s}^{-1}$  north-east wind. The wide southern intertidal banks (Poutawa and Hikihi) are the only other areas where pronounced wind-driven flows occur, up to  $0.1 \text{ m s}^{-1}$  downwind.

For prevailing **south-west** winds (Fig. 9C,D), the wind-induced circulation exhibits similar magnitudes, but reversed, to the pattern which occurs during equivalent-speed north-east winds. This implies a near-linear response of the Manukau Harbour basin to wind-forcing of equal magnitude from the two dominant wind directions. During south-westerlies, a coherent north to north-east

wind-induced flow is established at higher tide levels along the north-eastern perimeter through to Mangere Inlet. This mechanism partly explains the large influx of suspended sediment measured by Green & Bell (1995) on the southern side of Mangere Inlet during the late stages of a flood tide, after a persistent south-west breeze of  $6\text{--}8 \text{ m s}^{-1}$ , together with locally-sourced sediment stirred from the adjacent nearshore surf zone. Wind-induced flows (depth-averaged) in the main channels are upwind, which is to the south-west in the instance of lower Purakau Channel, virtually cancelling out the tidal residual current.

Thus, the wind-driven circulation in Manukau Harbour is characterised by downwind-directed flows on the shallow intertidal sand banks with pressure-driven return flows (vertically-averaged) in the deeper main channels of the inner harbour.

This same pattern was demonstrated by Glorioso & Davies (1995) using a three-dimensional model for an estuary with a deep central channel flanked by shallow areas. In these instances for well-mixed estuaries, the circulation and flushing time are determined largely by horizontal processes and, hence, are fairly insensitive to the magnitude of the vertical eddy viscosity. As expected, the wind effect on depth-averaged flows in the deep entrance channel is negligible, although surface currents may be significantly affected by strong winds.

### Bottom stress, dissipation rate and tidal stirring

The mean bottom frictional stress ( $\tau_\beta$ ) during a tidal cycle  $T$  (Bowman et al. 1980), which is of interest in sediment transport, is given by:

$$\tau_\beta = \frac{1}{T} \int_0^T \frac{\rho g}{C^2} |\tilde{U}|^2 dt \quad (2)$$

where  $\rho$  = density of sea water, assumed constant;  $g$  = acceleration as a result of gravity; and  $\tilde{U}$  is the depth-averaged current velocity. As expected from the current velocities, the regions of highest bottom stress are located in the entrance (off Ninepin Rock) and along the southern side of the entrance channel seaward of Te Pirau Point, with a maximum value of  $3.6 \text{ N m}^{-2}$  in the entrance. Mean bottom stresses above  $0.1 \text{ N m}^{-2}$  occur in the western half of the harbour (seaward of the  $174^\circ 41' \text{ E}$  meridian), plus the entrance to Waiuku River tidal inlet. Low tidal bottom stresses ( $< 0.01 \text{ N m}^{-2}$ ) occur over the intertidal flats arising from both relatively low velocities and shorter inundation periods, but conversely bottom stresses generated by wind-waves are much higher in these shallow areas causing high turbidities (Bell et al. 1997; Green et al. 1997). The dissipation rate of  $M_2$  energy by bottom friction, which represents the rate of work per unit area made by the water column against bottom friction (where the current velocity in Equation 2 is cubed; Bowman et al. 1980), is similar to the bottom stress distribution. High dissipation rates ( $> 0.05 \text{ W m}^{-2}$ ) occur in the entrance channel (maximum off Ninepin Rock of  $5.3 \text{ W m}^{-2}$ ) and the seaward ends of the inner-harbour channels. As an indication of the ability of vertical tidal stirring to inhibit summer stratification, Simpson & Hunter (1974) derived a stratification index ( $h/u^3$ ), where at any location,  $h$  is the water depth, and  $u$  is the amplitude of the mean-spring tidal current. Low values indicate areas which remain well mixed

vertically. Intermediate values in the range  $30\text{--}100$  up to  $500 \text{ m}^{-2} \text{ s}^3$  (Simpson & Hunter 1974; Bowman et al. 1980; Argote et al. 1995) have been associated with frontal structures found in shelf seas and produced by variations in tidal mixing. Values of this index in the entrance to Manukau Harbour are low ( $< 5 \text{ m}^{-2} \text{ s}^3$ ) as expected, because of high tidal dissipation rates. For the main channels dissecting the inner Harbour, the stratification index is mainly in the range  $20\text{--}35 \text{ m}^{-2} \text{ s}^3$ . Therefore summer stratification is unlikely, as confirmed by regular water-column observations.

### CONCLUSIONS

A two-dimensional finite-element model, which handles flooding/drying of intertidal areas, complemented by field data, has produced information on tidal propagation in Manukau Harbour with reasonably good accuracy. Excluding the upper-harbour inlets, c. 50% of both the amplification and phase lag in the dominant  $M_2$  constituent occurs through the deep 10-km-long entrance channel to Puponga Point. The calculation of tidal residuals has isolated flood-directed and ebb-directed residual circulation respectively on either side of Puponga Point, which is a typical pattern for alternating flows around headlands. The model has established the dominance of tidal-driven over wind-driven circulation in most of the harbour except higher up on the intertidal areas. Wind-driven circulation in the inner harbour is characterised by downwind flows over intertidal sand banks with pressure-driven return flows (vertically-averaged) in the deeper main channels. Tidal dissipation rates are sufficiently high to inhibit the onset of any summer stratification.

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## APPENDIX 1

### Model calibration and performance

#### Calibration and verification

Calibration of the model was achieved by simulating the dominant  $M_2$  tide where the open-boundary water elevation at time  $t$  was computed as:

$$\eta(t) = \sum_i f_i A_i \cos(\omega_i t - g_i + E_i) \quad (3)$$

where  $A_i$ ,  $g_i$ ,  $\omega_i$ ,  $f_i$ , and  $E_i$  are respectively the amplitude, phase lag, angular frequency, nodal modulation factor and equilibrium argument of the  $i$ th tidal constituent, in this instance the  $M_2$  (Walters 1987). The few historic measurements available over a wide swath along the north-western coast (Heath et al. 1977) all indicated that  $M_2$  phases are clustered around  $289^\circ$ . The nearest node in the CSU world ocean tidal model (Egbert et al. 1994), 45 km offshore from Manukau Harbour, yielded an  $M_2$  amplitude and phase of 109 cm and  $268^\circ$  (relative to NZST). To match the observed  $M_2$  tide at the harbour entrance (Paratutae Island), the optimum values of  $M_2$  amplitude and phase on the open-sea boundary (Fig. 2) were found to be 110 cm and  $276.5^\circ$  (or 4.5 ahead of Paratutae), which compare favourably with the CSU model results. The spatial distribution of the friction coefficient (Chezy  $C$ ), which yielded the best overall match with water level and current velocity field data, was  $C = 64 \text{ m}^{0.5} \text{ s}^{-1}$  over most of Manukau Harbour, except for the Papakura Channel, the ebb-tide delta and offshore waters beyond the entrance, where  $C = 54 \text{ m}^{0.5} \text{ s}^{-1}$ . These Chezy values are within the range of other

model studies (e.g., Dronkers 1964). The primary calibration of the tidal model was based on matching  $M_2$  tide heights at eight harbour sites (Table 4). The comparison between observed and model tide elevations for Onehunga Wharf is shown in Fig. 10A (the most distant site from the entrance). The secondary calibration step involved comparing the simulated  $M_2$  velocity components with available field data, mainly Papakura Channel (Sites 4, 5) and Mangere Inlet near Onehunga Wharf (Site 8). At these sites, model peak-ebb speeds were different to measured values by  $-8$ ,  $-14$ , and  $-12\%$  respectively, whereas for model peak-flood speeds, the differences were  $-10$ ,  $+13$ , and  $-12\%$ . The differences in peak-velocity phases were within 0.5 h. Although model velocities at these sites were mostly lower than measured values, the overall match is reasonable considering that the single-point current measurements have been scaled to an equivalent velocity at 0.4 of the depth above the bed and then compared to depth-averaged results from the nearest mesh node in the model.

The spring-tide model run, was used to verify the performance of the calibrated tidal-model, by freezing all model parameters such as the friction and eddy viscosity coefficients. For the offshore open boundary conditions, actual spring-tide water levels from the harbour entrance were transformed by the same phase difference and amplitude factor used in the calibration runs. The resulting match between the observed and

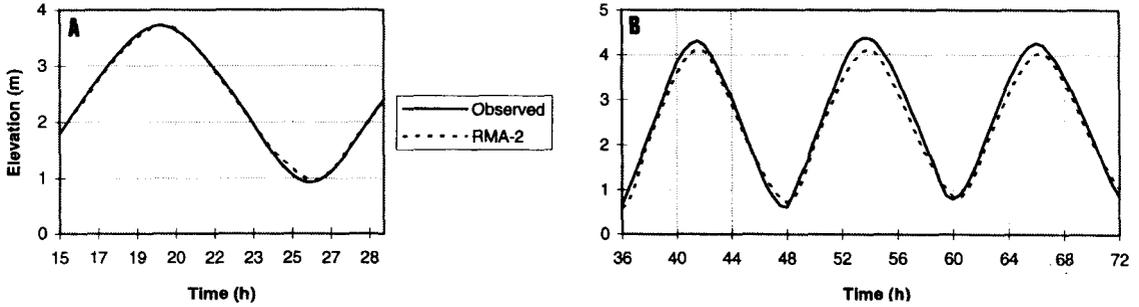


Fig. 10 Water level time series at Onehunga Wharf, New Zealand, comparing observed with modelled data for the: **A**,  $M_2$  calibration run; and **B**, spring-tide verification run.

model water levels at Onehunga Wharf is shown in Fig. 10B. The simulated water level around high water for the middle cycle is somewhat lower ( $-0.19$  m) than the observed spring high water at 53.5 h and similarly on the following ebb tide (partly because of set-up caused by a south-westerly wind peaking at  $14$  m  $s^{-1}$ ), but otherwise the match at the low waters and the tide phasing is reasonable. A closer match was obtained at both other sites (Sites 2, 5) where concurrent observations were available. The comparisons between observed and simulated tidal elevations and velocities are reasonable enough to conclude that the tidal model produces reliable results.

#### *Tidal model evaluation*

The tidal model was evaluated mainly by comparing  $M_2$  model results with observations from eight sites around Manukau Harbour, which were mainly short-term deployments of c. 1 month (Table 4). Discrepancies between model results and field observations are the result of both uncertainties in the observed harmonics (Foreman & Walters 1990; Foreman et al. 1995) and inaccuracies in the model formulation and application.

Most of the observed records were no more than 1 month, so to reduce inaccuracies data measured during concurrent time periods was used wherever possible to optimise consistency. Inaccuracies are also present with the model, particularly in the way it is applied, such as the schematisation of the basin bathymetry and friction factor and in the instance of the  $M_2$  simulation, interactions via the non-linear friction term with other Group 1 and 2 constituents are neglected. In particular for the Manukau Harbour, further refinements of the spatial distribution of the bed-friction coefficient and the finite-element mesh may be needed in the future, especially to redress the paucity of bathymetric data on the intertidal banks and southern inlets and improve resolution in the narrower channels. This should be complemented with longer term ( $> 1$  month) field observations, to improve the match between observed and model simulations. Provided the finite-element mesh is carefully designed by aligning element boundaries with isobaths on intertidal areas, the application of the RMA-2 model described in this paper has demonstrated that finite-element models are capable of reliable simulation of estuarine hydrodynamics in harbours with extensive intertidal sandflats.